

Lecture 6 : Analysis techniques

Decomposing the circulation

We are interested in the time- and zonal-mean circulation, the deviations from this mean, and how the two interact.

Define time mean

$$\bar{A} = \frac{1}{\tau} \int_0^{\tau} A dt$$

and deviations: $A' = A - \bar{A}$

$$\text{note that } \overline{A'} = \overline{A - \bar{A}} = \bar{A} - \bar{A} = 0$$

τ is usually seasonal or longer timescale

Also define zonal mean

$$[A] = \frac{1}{2\pi} \int_0^{2\pi} A d\lambda$$

$$A^* = A - [A]$$

$$[A^*] = 0$$

Why do this?

Useful for analysing covariances

$$AB = (\bar{A} + A')(\bar{B} + B')$$

$$= \bar{A}\bar{B} + \bar{A}B' + \bar{B}A' + A'B'$$

Take the time mean:

$$\overline{AB} = \overline{\overline{AB}} + \overline{\overline{A'B'}} + \overline{\cancel{BA'}} + \overline{A'B'}$$

$$= \overline{AB} = \overline{A'B'} = 0$$

$$\overline{AB} = \overline{AB} + \overline{A'B'}$$

↑
covariance!

Example I

Consider the conservation equation for scalar χ

$$\frac{D\chi}{Dt} = 0$$

Write this in flux form.

Use tangent plane equations in pressure coordinates for simplicity:

$$\frac{\partial \chi}{\partial t} + \underline{u} \cdot \nabla_p \chi = \frac{\partial \chi}{\partial t} + u \frac{\partial \chi}{\partial x} + v \frac{\partial \chi}{\partial y} + \omega \frac{\partial \chi}{\partial p} = 0 \quad (1)$$

Also have mass continuity:

$$\nabla_p \cdot \underline{u} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial \omega}{\partial p} = 0 \quad (2)$$

Add (1) & $\chi \times (2)$:

$$\frac{\partial \chi}{\partial t} + \underline{u} \cdot \nabla_p \chi + \chi \nabla_p \cdot \underline{u} = 0$$

$$\frac{\partial \chi}{\partial t} + \nabla_p \cdot (\underline{u} \chi) = 0 \quad \Rightarrow \text{Flux form}$$

Now take zonal mean:

$$\frac{\partial \langle \chi \rangle}{\partial t} + \nabla_p \cdot \langle \underline{u} \chi \rangle = 0$$

$$\text{also } \nabla_p \cdot \langle \underline{u} \rangle = \frac{\partial \langle u \rangle}{\partial z} + \frac{\partial \langle \omega \rangle}{\partial p} = 0 \quad \left[\lambda\text{-derivative is zero} \right]$$

Now use our decomposition:

$$\langle \underline{u} \chi \rangle = \langle \underline{u} \rangle \langle \chi \rangle + \langle \underline{u}^* \chi^* \rangle$$

$$\frac{\partial \langle \chi \rangle}{\partial t} + \nabla_p \cdot \langle \underline{u} \rangle \langle \chi \rangle = - \nabla_p \cdot \langle \underline{u}^* \chi^* \rangle$$

$$\Rightarrow \frac{\partial \langle \chi \rangle}{\partial t} + \langle \underline{u} \rangle \cdot \nabla_p \langle \chi \rangle + \langle \chi \rangle \nabla_p \cdot \langle \underline{u} \rangle = - \nabla_p \cdot \langle \underline{u}^* \chi^* \rangle$$

by zonal-mean continuity

$$\boxed{\frac{\partial \langle \chi \rangle}{\partial t} + \langle \underline{u} \rangle \cdot \nabla_p \langle \chi \rangle = - \frac{\partial \langle \underline{u}^* \chi^* \rangle}{\partial y} - \frac{\partial \langle \omega^* \chi^* \rangle}{\partial p}} \quad \left| \begin{array}{l} \lambda\text{-derivative} \\ \text{of divergence} \\ \text{is zero} \end{array} \right.$$

Can do this derivation in spherical coordinates

Remember $\frac{D\chi}{Dt} = \frac{\partial \chi}{\partial t} + \underline{u} \cdot \nabla_p \chi = \frac{\partial \chi}{\partial t} + \frac{u}{R_e \cos \phi} \frac{\partial \chi}{\partial \lambda} + \frac{v}{R_e} \frac{\partial \chi}{\partial \phi} + \omega \frac{\partial \chi}{\partial p}$

$$\Rightarrow \frac{\partial \chi}{\partial t} + \frac{u}{R_e \cos \phi} \frac{\partial \chi}{\partial \lambda} + \frac{v}{R_e} \frac{\partial \chi}{\partial \phi} + \omega \frac{\partial \chi}{\partial p} = 0 \quad (1)$$

Also, from continuity: $\frac{1}{R_e \cos \phi} \left[\frac{\partial u}{\partial \lambda} + \frac{\partial v \cos \phi}{\partial \phi} \right] + \frac{\partial \omega}{\partial p} = 0 \quad (2)$

Take (1) + $\chi \times$ (2):

$$\frac{\partial \chi}{\partial t} + \frac{u}{R_e \cos \phi} \frac{\partial \chi}{\partial \lambda} + \frac{\chi}{R_e \cos \phi} \frac{\partial u}{\partial \lambda} + \frac{v}{R_e} \frac{\partial \chi}{\partial \phi} + \frac{\chi}{R_e \cos \phi} \frac{\partial v \cos \phi}{\partial \phi} + \chi \frac{\partial \omega}{\partial p} + \omega \frac{\partial \chi}{\partial p} = 0$$

$$\frac{\partial \chi}{\partial t} + \frac{1}{R_e \cos \phi} \left[\frac{\partial (\chi u)}{\partial \lambda} \right] + \frac{1}{R_e \cos \phi} \left[\frac{\partial \chi v \cos \phi}{\partial \phi} \right] + \frac{\partial (\chi \omega)}{\partial p} = 0$$

$$\frac{\partial \chi}{\partial t} + \frac{1}{R_e \cos \phi} \left[\frac{\partial (u \chi)}{\partial \lambda} + \frac{\partial (v \chi \cos \phi)}{\partial \phi} \right] + \frac{\partial (\omega \chi)}{\partial p} = 0$$

$$\frac{\partial \chi}{\partial t} + \nabla_p \cdot (\underline{u} \chi) = 0$$

Take zonal mean

$$\frac{\partial [\chi]}{\partial t} + \frac{1}{R_e \cos \phi} \frac{\partial [v \chi] \cos \phi}{\partial \phi} + \frac{\partial [\omega \chi]}{\partial p}$$

The λ -derivative of the divergence is zero.

Remember $[v \chi] = [v][\chi] + [v^* \chi^*]$

$$\Rightarrow \frac{\partial [\chi]}{\partial t} + \frac{1}{R_e \cos \phi} \frac{\partial [v][\chi] \cos \phi}{\partial \phi} + \frac{\partial [\omega][\chi]}{\partial p} = -\frac{1}{R_e \cos \phi} \frac{\partial [v^* \chi^*] \cos \phi}{\partial \phi} - \frac{\partial [\omega^* \chi^*]}{\partial p}$$

Expand LHS:

$$\begin{aligned} \frac{\partial [\chi]}{\partial t} + \frac{[v]}{R_e} \frac{\partial [\chi]}{\partial \phi} + [\omega] \frac{\partial [\chi]}{\partial p} &= -\frac{1}{R_e \cos \phi} \frac{\partial (v^* \chi^*) \cos \phi}{\partial \phi} - \frac{\partial (\omega^* \chi^*)}{\partial p} \\ + \frac{[\chi]}{R_e \cos \phi} \frac{\partial [v] \cos \phi}{\partial \phi} + [\chi] \frac{\partial [\omega]}{\partial p} & \end{aligned}$$

Now, the zonal-mean continuity equation may be written,

$$\nabla_p \cdot [\underline{u}] = 0$$

$$\Rightarrow \frac{1}{R_e \cos \phi} \frac{\partial [v] \cos \phi}{\partial \phi} + \frac{\partial [\omega]}{\partial p} = 0$$

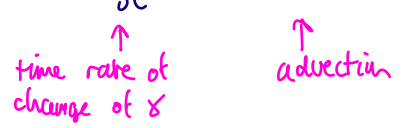
Hence the equation for the zonal-mean χ may be written:

$$\frac{\partial [\chi]}{\partial t} + \frac{[v]}{R_e} \frac{\partial [\chi]}{\partial \phi} + [\omega] \frac{\partial [\chi]}{\partial p} = -\frac{1}{R_e \cos \phi} \frac{\partial [v^* \chi^*] \cos \phi}{\partial \phi} - \frac{\partial [\omega^* \chi^*]}{\partial p}$$

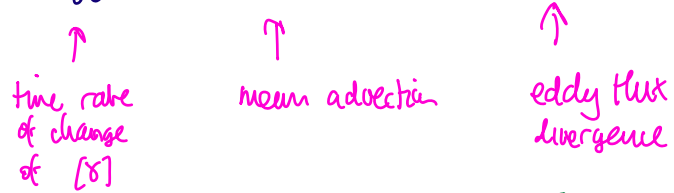
$$\frac{\partial [\chi]}{\partial t} + [\underline{u}] \cdot \nabla_p [\chi] = -\nabla_p \cdot [\underline{u}^* \chi^*]$$

Equation for the zonal mean:

$$\frac{\partial \bar{\chi}}{\partial t} + \underline{u} \cdot \nabla_p \bar{\chi} = 0$$



$$\frac{\partial [\bar{\chi}]}{\partial t} + [\underline{u}] \cdot \nabla_p [\bar{\chi}] = - \nabla_p \cdot [\underline{u}^* \chi^*]$$



 This is especially important to understand!

Example II

time	q (kg/kg)	U (ms ⁻¹)	ρ (kgm ⁻³)
1	0.01	3	1
2	0.002	-1	1
3	0.014	6	1

1. Calculate meridional flux of humidity

$$\text{Flux} = \rho_0 \overline{Uq}$$

$$\left\{ \begin{array}{l} \text{units: } \text{ms}^{-1} \cdot \text{kg/kg} \cdot \text{kg m}^{-3} \\ = \text{kg m}^{-2} \text{s}^{-1} \end{array} \right.$$

more accurately a flux density
(units in p-coords are more confusing)

$$= \left[0.01(3) + 0.002(-1) + 0.014(6) \right] \frac{1}{3}$$
$$= 0.037 \text{ kg m}^{-2} \text{ s}^{-1}$$

Flux by mean circulation

$$F_{\text{MEAN}} = \rho_0 \overline{Uq} = \rho_0 \left(\frac{3+6-1}{3} \right) \left(\frac{0.01+0.002+0.014}{3} \right)$$
$$= 0.023 \text{ kg m}^{-2} \text{ s}^{-1}$$

Flux by eddies

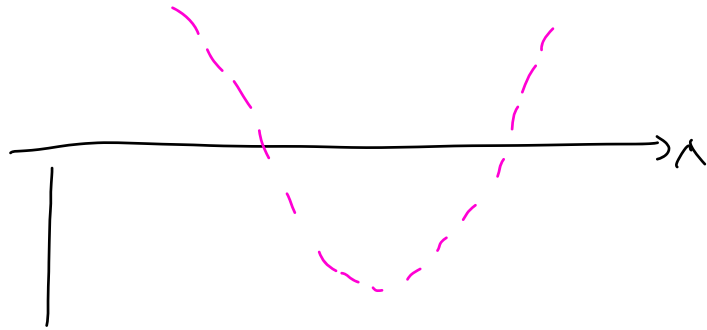
$$F_{\text{EDDY}} = \rho_0 \overline{U'q'} = \rho_0 \overline{Uq} - \rho_0 \overline{Uq} = 0.014 \text{ kg m}^{-2} \text{ s}^{-1}$$

Example III

$$[F] = \rho_0 [Uq]$$

$$U = \Delta V \cos \lambda$$
$$q = q_0 + \Delta q \cos \lambda$$

U, q



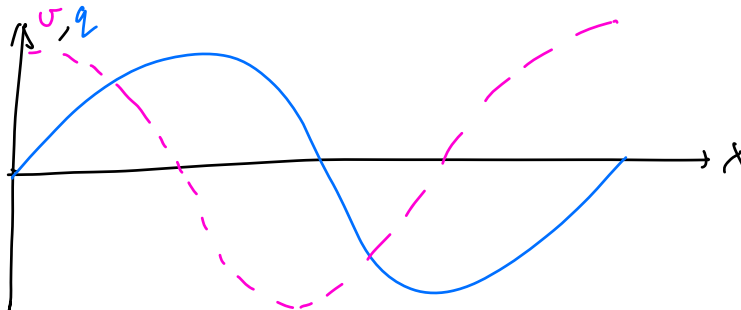
Note that $\bar{v} = 0$, $\bar{q} = 0$

$$[F] = \frac{\rho_0 \Delta q \Delta v}{2\pi} \int_0^{2\pi} \cos(\lambda) \cos(\lambda) d\lambda = \frac{\rho_0 \Delta q \Delta v}{2\pi} \int_0^{2\pi} \cos^2(\lambda) d\lambda > 0$$

$$\frac{1}{2\pi} \int_0^{2\pi} \cos^2(\lambda) d\lambda = \frac{1}{2\pi} \int_0^{2\pi} \frac{1 + \cos(2\lambda)}{2} d\lambda = \frac{1}{2}$$

$$\Rightarrow [F] = \frac{\rho_0 \Delta q \Delta v}{2}$$

now suppose $q = \sin \lambda$



$$[F] = \frac{\rho_0 \Delta v \Delta q}{2} \int_0^{2\pi} \cos(\lambda) \sin(\lambda) d\lambda$$

Let $u = \sin(\lambda)$

$$[F] = \frac{1}{2\pi} \int u du = \frac{1}{2\pi} \left[\frac{\sin^2 \lambda}{2} \right]_0^{2\pi} = 0$$

Combining the space & time operator

Consider the temporal covariance: \overline{AB}

We know that

$$\overline{AB} = \overline{A\overline{B}} + \overline{A'B'}$$

Now, take the zonal mean

$$[\overline{AB}] = [\overline{A\overline{B}}] + [\overline{A'B'}]$$

Decompose the first term into zonal mean and eddy:

$$[\overline{AB}] = [\overline{A}][\overline{B}] + [\overline{A^*B^*}] + [\overline{A'B'}]$$

zonal- and time-mean *stationary eddies* *transient eddies*

Could also go further:

$$[\overline{AB}] = [\overline{A}][\overline{B}] + [\overline{A^*B^*}] + [\overline{A'}][\overline{B'}] + [\overline{A'^*B'^*}]$$

But usually do not separate transient asymmetric from transient symmetric.

Example

Take $A=u$, $B=u$

$$[\overline{u^2}] = [\overline{u}]^2 + [\overline{u'^2}] + [\overline{u'^2}]$$

Add to $[\bar{v}^2]$:

$$[\overline{u^2 + v^2}] = [\bar{u}^2 + \bar{v}^2] + [\overline{u'^2}] + [\overline{v'^2}] + [\overline{u'v'}] + [\overline{v'u'}]$$

↑
mean kinetic energy

↑
KE of stationary eddies

↑
KE of transient eddies

State estimation techniques

How do we calculate $[\overline{u'^2}]$ or $[\overline{u'v'}]$ for the atmosphere?

Observations are scattered in space and time

Need a method to construct gridded estimate

Method 1: Peixoto & Oort - objective analysis

- Used to produce figures in Peixoto & Oort (1992).
- uses radiosonde data.

Step 1

- Quality control data (remove outliers)
- Calculate means + covariances at each station

$$\bar{T}, \bar{T'^2}, \bar{vT}, \bar{v'T'}$$

Step 2

- Define an "initial guess" on a grid

- Peixoto & Oort used a zonal mean climatology in 10° latitude bands

step 3

- Update the initial guess using the available data.
- In data sparse regions solve the following BVP:

$$\nabla^2 T = F(\lambda, \phi)$$

$$F = \nabla^2 T_{\text{guess}}$$

Why not just a simple interpolation routine?

Method 2: Modern reanalysis

- Use atmospheric general circulation model to produce initial guess
- Correct initial guess using data assimilation techniques
- Basic example is the Kalman Filter
 - initial guess provided by forecast model
 - updated based on observations in order to minimise a cost function that is constructed based on estimated error in obs and model.

Next lecture: Begin to apply these techniques to understand the atmosphere's angular momentum budget.